

**Readings:** Woodhouse - General Relativity  
                   Chapter 6 - Einstein's Equation  
                   Chapter 7 - Spherical Symmetry  
                   Walecka - General Relativity  
                   Chapter 7 - General Relativity

**Present Ideas and Lead Discussion:**

Woodhouse Chapter 6	Erin_____
Woodhouse Chapter 7	Ben P_____
Walecka Chapter 7	Robert_____

**Present Problems:**

Wo7.1 - Time	Erin_____
Wo7.2 - Lie Derivative	Ben P_____
Wo7.3 - More Lie derivatives	
Wa7.1 - How many independent components?	Robert_____
Wa7.2 - Ricci=0 does not imply Riemann=0	
Wa7.3 - Manipulating geodesic equations	Emma_____
Wa7.4 - Newtonian limit	
Wa7.5 - Circular motion	Eric_____
Wa7.6 - Radial motion	
Wa7.8 - Time difference	Markus_____
Wa7.9 - Length of meter stick	

**EP27. Timelike Geodesics**

Chris\_\_\_\_\_

Find the timelike geodesics for the metric

$$ds^2 = \frac{1}{t^2}(-dt^2 + dx^2)$$

**EP28. More Geodesics**

Chris\_\_\_\_\_

Consider the 2-dimensional metric

$$ds^2 = a^2(d\chi^2 + \sinh^2 \chi d\phi^2)$$

- (a) Compute the connection coefficients  $\Gamma_{jk}^i$
- (b) Compute all components of the Riemann tensor,  $R_{kl}^{ij}$ , the Ricci tensor,  $R_j^i$ , and the Ricci scalar  $R$ .
- (c) A geodesic starts at  $\chi=b$ ,  $\phi=0$  with tangent  $d\phi/d\lambda=1$ ,  $d\chi/d\lambda=0$ . Find the trajectory  $\chi(\phi)$ .
- (d) A second geodesic starts at  $\chi=b+\xi$  ( $\xi \ll 1$ ), also initially in

the  $\phi$ -direction. How does the separation initially increase or decrease along the two curves.

- (e) What is the shape of the geodesic trajectory as  $a \rightarrow \infty$ ,  $\chi \rightarrow 0$  with  $r = a\chi$  fixed.

**EP29. Parallel Transport on a Sphere (Grade #2)** \_Everyone\_

- (a) On the surface of a 2-sphere of radius  $a$ ,

$$ds^2 = a^2(d\theta^2 + \sin^2\theta d\phi^2)$$

Consider the vector  $\vec{A}_0 = \vec{e}_\theta$  at  $\theta = \theta_0$ ,  $\phi = 0$ . The vector is parallel transported all the way around the latitude circle  $\theta = \theta_0$  (i.e., over the range  $0 \leq \phi \leq 2\pi$  at  $\theta = \theta_0$ ). What is the resulting vector  $\vec{A}$ ? What is its magnitude  $(\vec{A} \cdot \vec{A})^{1/2}$ ?

HINT: derive differential equations for  $A^\theta$  and  $A^\phi$  as functions of  $\phi$ .

**EP30. Curvature on a Sphere** \_Ben G\_

- (a) Compute all the nonvanishing components of the Riemann tensor  $R_{ijkl}$  ( $(i,j,k,l) \in (\theta,\phi)$ ) for the surface of a 2-sphere.
- (b) Consider the parallel transport of a tangent vector  $\vec{A} = A^\theta \vec{e}_\theta + A^\phi \vec{e}_\phi$  on the sphere around an infinitesimal parallelogram of sides  $\vec{e}_\theta d\theta$  and  $\vec{e}_\phi d\phi$ . Using the results of part (a), show that to first order in  $d\Omega = \sin\theta d\theta d\phi$ , the length of  $\vec{A}$  is unchanged, but its direction rotates through an angle equal to  $d\Omega$ .
- (c) Show that, if  $\vec{A}$  is parallel transported around the boundary of any simply connected solid angle  $\Omega$ , its direction rotates through an angle  $\Omega$ . ("Simply connected" is a topological term meaning that the boundary of the region could be shrunk to a point; it tells us that there are no holes in the manifold or other pathologies). Using the result of part (b) and intuition from proofs of Stoke's theorem, this should be an easy calculation. Compare with the result of **EP30**.

**EP31. Riemann Tensor for 1+1 Spacetimes** \_Sam\_

- (a) Compute all the nonvanishing components of the Riemann tensor for the spacetime with line element

$$ds^2 = -e^{2\phi(x)} dt^2 + e^{2\psi(x)} dx^2$$

- (b) For the case  $\phi = \psi = \frac{1}{2} \ln |g(x - x_0)|$  where  $g$  and  $x_0$  are constants, show that the spacetime is flat and find a coordinate

transformation to globally flat coordinates  $(\bar{t}, \bar{x})$  such that  
 $ds^2 = -d\bar{t}^2 + d\bar{x}^2$ .

**EP32. About Vectors Tangent to Geodesics**

\_Margaret\_

Let  $x^\mu(\tau)$  represent a timelike geodesic curve in spacetime, where  $\tau$  is the proper time as measured along the curve. Then  $u^\mu \equiv dx^\mu / d\tau$  is tangent to the geodesic curve at any point along the curve.

(a) If  $g_{\mu\nu}$  is the metric of spacetime, compute the magnitude of

the vector  $u^\mu$ . Do not use units where  $c = 1$ , but keep any factors of  $c$  explicit. Compare your result with the one obtained in flat Minkowski spacetime. HINT: The magnitude of a timelike vector  $v^\mu$  is given by  $(-g_{\mu\nu}v^\mu v^\nu)^{1/2}$ .

(b) Consider a contravariant timelike vector  $v^\mu$  at a point P on the geodesic curve. Move the vector  $v^\mu$  from the point P to an arbitrary point Q on the geodesic curve via parallel transport. Prove that the magnitude of the vector  $v^\mu$  at the point Q equals the magnitude of the vector  $v^\mu$  at point P.

(c) Suppose that at the point P on the geodesic curve,  $v^\mu = u^\mu$ . Now, parallel transport the vector  $v^\mu$  along the geodesic curve to arbitrary point Q. Show that  $v^\mu = u^\mu$  at the point Q.

NOTE: This result implies that a vector tangent to a geodesic at a given point will always remain tangent to the geodesic curve when parallel transported along the geodesic.